

A Biophysically Based Mathematical Model of Unitary Potential Activity in Interstitial Cells of Cajal

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ABSTRACT Unitary potential (UP) depolarizations are the basic intracellular events responsible for pacemaker activity in interstitial cells of Cajal (ICCs), and are generated at intracellular sites termed “pacemaker units”. In this study, we present a mathematical model of the transmembrane ion flows and intracellular Ca^{2+} dynamics from a single ICC pacemaker unit acting at near-resting membrane potential. This model quantitatively formalizes the framework of a novel ICC pacemaking mechanism that has recently been proposed. Model simulations produce spontaneously rhythmic UP depolarizations with an amplitude of ~ 3 mV at a frequency of 0.05 Hz. The model predicts that the main inward currents, carried by a Ca^{2+} -inhibited nonselective cation conductance, are activated by depletion of sub-plasma-membrane $[\text{Ca}^{2+}]$ caused by sarcoplasmic reticulum calcium ATPase Ca^{2+} sequestration. Furthermore, pacemaker activity predicted by our model persists under simulated voltage clamp and is independent of $[\text{IP}_3]$ oscillations. The model presented here provides a basis to quantitatively analyze UP depolarizations and the biophysical mechanisms underlying their production.

INTRODUCTION

Phasic gastrointestinal (GI) muscle tissue generates rhythmic contractions without input from the enteric nervous system (1). The electrical activity that coordinates these contractions is termed “slow waves.” Originally thought to be an intrinsic property of smooth muscle cells (SMCs) (2), slow waves are now known to originate from a specialized group of pacemaker cells called interstitial cells of Cajal (ICCs). This conclusion was based in part on experimental investigations wherein GI tissue lacking an intact ICC network failed to produce slow-wave electrical activity (3–5).

Slow waves generated by ICCs, also referred to in the literature as pacemaker potentials (PP), conduct via gap junctions to SMC, activating L-type Ca^{2+} channels, Ca^{2+} influx, and contraction of muscle cells (4–6). Full slow waves result from the summation of a large number of small-amplitude and localized cellular membrane fluctuations, known as unitary potentials (UPs), which have been detected in many GI muscles (7–10). UPs are thus the elementary events responsible for ICC pacemaker activity, and hence are fundamental in driving coordinated GI muscle contractions.

Despite the importance of UPs in generating ICC pacemaker activity, controversy remains as to the biophysical mechanisms and processes underlying their production (see discussion in Sanders et al. (11)). Studies on intact muscle strips, in which the activity of ICCs cannot be isolated, have suggested a variety of mechanisms for pacemaker activity (7,12–14). Pharmacological and ultrastructural experimental

studies performed on ICCs over the last two decades have led to the formation of a novel pacemaking hypothesis (11,15). This hypothesis, which we will refer to henceforth as the Sanders pacemaking hypothesis, proposes a detailed structural feature, occurring at specific sites within ICCs (termed pacemaker units), which provides a unique intracellular Ca^{2+} -handling mechanism that appears to be fundamental to generation of spontaneous pacemaker currents.

In brief, the hypothesis states that the pacemaker cycle is initiated by release of Ca^{2+} from inositol-1,4,5-triphosphate-receptor (IP_3R)-operated stores that are in close proximity to the plasma membrane and to mitochondria. The increase in local $[\text{Ca}^{2+}]$ gates opens a mitochondrial Ca^{2+} uptake mechanism that results in a transient reduction in $[\text{Ca}^{2+}]$ in the cytoplasmic microdomain of the pacemaker unit. The reduction in sub-plasma-membrane $[\text{Ca}^{2+}]$ activates the primary pacemaker conductance that experiments on isolated ICCs have shown to be a Ca^{2+} -inhibited, nonselective cation conductance. The pacemaker conductance is voltage-independent, so a voltage-dependent mechanism is required to coordinate the many pacemaker units within a single ICC and within a network of ICCs. This voltage-dependent mechanism is likely to be a T-type Ca^{2+} conductance expressed by ICCs (16,17), which is activated by the transient depolarizations of UPs and provides a means of nonregenerative propagation of slow-wave potentials over many millimeters. For a full description of the Sanders pacemaker hypothesis, and the experimental studies that led to its formation, the reader is directed to a recent review (11).

Due to the small dimensions of the pacemaker unit, however, it is technically difficult to empirically investigate the Sanders hypothesis in a direct manner. Therefore, the application of biophysically based mathematical models provides an alternative approach to test the underlying assumptions

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about the pacemaker cycle. The biophysically based modeling approach has been used successfully, in combination with physiological data, for the last 40 years to investigate cardiac cellular activity (18). Compared to the cardiac field, modeling of ICC pacemaker activity is still in its infancy. Nonetheless, several investigators have created mathematical models that reproduce specific aspects of ICC pacemaker activity (12,13,19,20). A brief summary of these established ICC models is given in Table 1. From this collection, note that only the model in Edwards and Hirst (12) has a UP basis to reproduce whole-cell electrical activity. However, the UP representation they use is phenomenological and therefore is unable to capture the subcellular Ca^{2+} -handling mechanisms thought to be involved in UP depolarizations.

In this study, we integrate experimental data and the understanding of the Sanders pacemaker hypothesis to introduce the first biophysically based computational modeling framework that simulates ICC pacemaker activity on the pacemaker unit spatial scale. This modeling framework is designed to reproduce UP depolarizations from a single pacemaker unit operating at resting membrane potential (RMP), which is approximately -70 mV in ICCs (17,21). Due to the complexity of the modeling framework, we introduce each of the model components separately as follows:

1. Parameter estimation
2. Intracellular compartmental volumes
3. Model state variables
4. Plasma membrane currents
5. The IP_3 receptor model
6. Other intracellular Ca^{2+} fluxes
7. Mitochondrial Ca^{2+} buffering

A schematic diagram of the pacemaker unit illustrating all compartmental volumes and ionic conductances, together with their respective interactions, is shown in Fig. 1.

MODEL FRAMEWORK

Parameter estimation

Where available, model parameters were either obtained by fitting to appropriate ICC experimental data or were based on values from the most appropriate alternative cell types, such as pancreatic acinar cells and sympathetic neurons. These

cells were deemed to be appropriate alternatives, because, like ICCs, they are nonmuscle cells and their activity is highly dependent on cycling of intracellular Ca^{2+} (22,23). All model parameters ascertained in this manner are specifically stated below. The remaining model parameters, for which there was no suitable data or appropriate alternative cell type, were optimized using a nonlinear fitting approach. This fitting procedure sought to optimize parameter values such that the model response reproduced experimentally observed characteristic UP features, namely,

- Amplitude, $Y_A = 3$ mV,
- Half-width, $Y_{HW} = 0.21$ s (defined as the time where membrane potential is $>50\%$ of Y_A),
- Frequency, $Y_F = 0.05$ Hz.

The values for Y_A and Y_{HW} stated above were obtained from experimental data recorded in guinea pig antral tissue (9). The exact frequency of UP depolarizations in ICCs is unknown, as it is not possible to isolate electrical activity from a single pacemaker unit. Therefore, we have assumed that Y_F is equal to the frequency of pacemaker potential depolarizations, which is ~ 3 cycles/min (or 0.05 Hz) in guinea pig antral tissue (9,21). It should be noted that although the modeling framework is based on ICC experimental data from gastric tissue, and hence is representative of UPs from this organ, the biophysical nature of the model allows for the potential applicability to GI-organ ICCs.

The optimization objective function, F , was calculated as the sum of the absolute residuals between the model and experimentally recorded characteristic values,

$$F = \bar{Y}_A + \bar{Y}_{HW} + \bar{Y}_F, \quad (1)$$

where $\bar{Y}_X = |Y_{X(\text{Experimental})} - Y_{X(\text{Model})}|$.

In addition to minimizing F , a further criterion was imposed on the parameters to ensure a physiological response from the model. This criterion was the enforcement of pacemaker activity cessation with the effective blocking of an individual ionic conductance, leaving the other conductances unaltered. This test was repeated for each ionic conductance within the pacemaker unit (Fig. 1). Effective blocking of an ionic conductance was simulated by setting the conductance rate to zero. To achieve a feasible parameter set that optimized F within this discrete constraint, we employed a Monte Carlo minimization algorithm (24).

Intracellular compartmental volumes

The intracellular compartmental volumes affect the concentrations of the major ionic species, and hence the pacemaker unit's ionic conductances. The spatial components used to represent the compartmental volumes within the model are based upon measurements from high-powered ICC electron micrographs (EM) (11,15). These EM studies clearly show

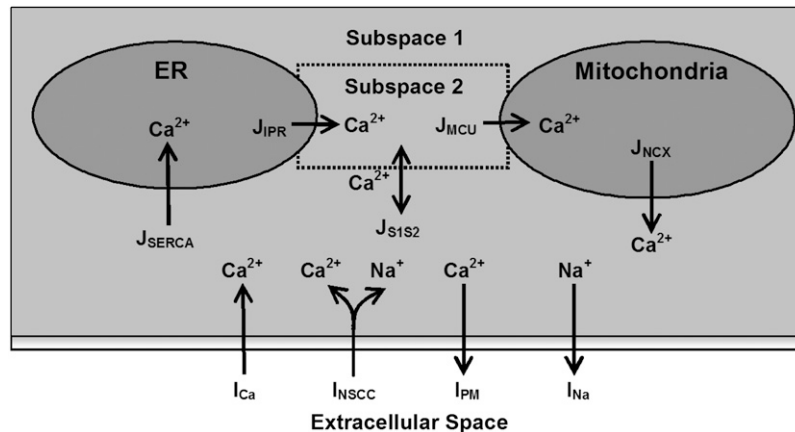
TABLE 1 Features of established ICC models

Reference	No. of state variables	Model type*	Behavior [†]	Tracks ionic species	UP basis	Spatial scale
19	2	P	D	No	No	SC/MC
12	5	P	S	No	Yes	MC
13	4	P	D	Yes	No	SC/MC
20	14	B	D	Yes	No	SC

SC, single cell; MC, multiple cells.

*P, phenomenological; B, biophysically based.

[†]D, deterministic; S, stochastic.



components included in this model are introduced in the Model framework section. The equations describing the ionic conductances and governing the model-state variables are located in Appendix A, and all model parameter and initial-state variable values are given in Table 2.

the location of an abundance of endoplasmic reticulum (ER) and mitochondria in close apposition to each other and the cellular membrane (Fig. 2). From these images, the total pacemaker unit volume was estimated to be on the order of $1 \mu\text{m}^3$.

Ionic concentrations within each compartmental volume are assumed to be homogeneous. This assumption is justified by the small volumes of the compartments and their relatively high ionic diffusion rates. A consequence of assuming ion concentration homogeneity within the compartmental volumes is that the model equations are only dependent upon time, and as such model dynamics can be described by ordinary differential equations. The spatial components of the pacemaker unit are modeled by four compartmental volumes: 1), the ER; 2), the mitochondria; 3), cytosolic subspace 1 (S_1); and 4), cytosolic subspace 2 (S_2).

Estimates of the intracellular compartmental volume ratios were made from ICC EMs (Fig. 2). Note that the volume

ratios are quoted relative to the smallest compartmental volume, S_2 , and not the pacemaker unit. By using this convention the number of free model parameters is reduced.

Endoplasmic reticulum subspace

The ER is an intracellular Ca^{2+} store that cycles Ca^{2+} by release via the IP_3R and is sequestered back again by the sarcoendoplasmic reticulum calcium ATPase (SERCA) pump proteins. It should be noted that the ryanodine receptors may be present in ICCs as well. However, experimental evidence suggests that ryanodine has little effect on pacemaker activity in ICCs (25) or in strips of guinea pig muscle (14). Therefore, a representation of ryanodine receptors was not included in our UP model.

Mitochondrial subspace

Mitochondrial Ca^{2+} uptake regulates cellular ATP production and defines the spatiotemporal transients of the cytosolic Ca^{2+} signals (26,27). Ca^{2+} is accumulated into the mitochondrial matrix by flux through the mitochondrial calcium uniporter (MCU) and released by the mitochondrial sodium/calcium exchanger (NCX).

Cytosolic subspaces

The remainder of the pacemaker unit consists of the cytosolic volume (CV). The close spatial arrangement of the ER and mitochondria gives rise to the possible formation of microdomains between these organelles. Such microdomains are consistent with a hypothesis known as the “hotspot theory” (28,29). In brief, the theory postulates that if the ER and mitochondria are spatially adjacent, then ER Ca^{2+} release can create regions of extremely high $[\text{Ca}^{2+}]$ between the organelles. Exposing the MCU to these Ca^{2+} “hotspots” activates this conductance, causing mitochondrial Ca^{2+}

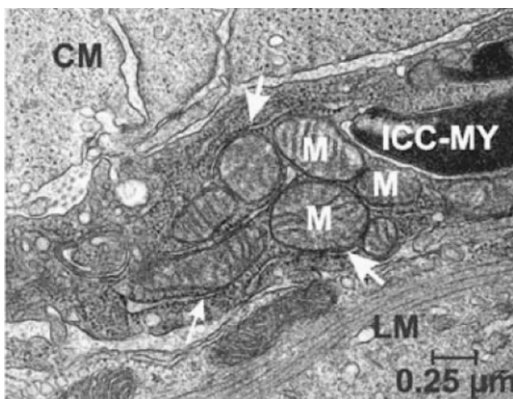


FIGURE 2 High-powered ICC electron micrograph photos showing the ultrastructural features. Note the abundance of mitochondria (M) and endoplasmic reticulum (arrowheads) in close proximity to each other and the cell membrane. Figure reprinted with permission from the Annual Review of Physiology, Volume 68 © 2006 by Annual Reviews (www.annualreviews.org).

accumulation. Such activation of the MCU would not be possible under normal cytosolic Ca^{2+} oscillations (0.1–1.0 μM), as the MCU is a low- Ca^{2+} -affinity conductance (30). Experimental evidence supporting this hypothesis has been observed in cells where there exists a close spatial relationship between the ER and mitochondria (27,28,31).

Due to the significantly higher Ca^{2+} fluxes through the aforementioned microdomains, a large $[\text{Ca}^{2+}]$ gradient will be formed between these regions and the rest of the CV. As a result, the assumption of ionic concentration homogeneity throughout the entire pacemaker unit CV is unrealistic. Therefore, the CV was split into two separate compartmental volumes, the bulk CV and a volume representing the microdomain formed between the organelles, termed cytosolic subspaces 1 (S_1) and 2 (S_2), respectively. S_1 regulates the Ca^{2+} fluxes into the ER and from the mitochondria. The concentrations of the S_1 ionic species are the ionic concentrations the plasma membrane currents are exposed to. The model assumes that there is no ion diffusion from S_1 to neighboring pacemaker units; hence the only mechanism for ion efflux from the pacemaker unit is across the cell membrane. S_2 accommodates ER Ca^{2+} release from the IP_3R and mitochondrial Ca^{2+} uptake by the MCU. Note that S_2 $[\text{Ca}^{2+}]$ regulates both the IP_3R and the MCU. Also note that, as there is no physical separation between the cytosolic subspaces, Ca^{2+} can diffuse between these volumes. We model this passive Ca^{2+} diffusion by the “flux”, J_{S1S2} (see Intracellular Ca^{2+} fluxes).

Model state variables

The major intra- and extracellular ion concentrations and cellular membrane potential are represented by dynamic state variables (SVs) that are governed by a series of coupled differential equations. In total, there are seven SVs that are tracked by the UP model (refer to the structures given schematically in Fig. 1), namely,

1. membrane potential (V_m),
2. S_1 $[\text{Ca}^{2+}]$ (C_{S1}),
3. S_1 $[\text{Na}^+]$ (N_{S1}),
4. S_2 $[\text{Ca}^{2+}]$ (C_{S2}),
5. ER $[\text{Ca}^{2+}]$ (C_{ER}),
6. mitochondrial $[\text{Ca}^{2+}]$ (C_{MT}), and
7. IP_3R control SV (H).

Plasma membrane conductances

The plasma membrane currents represent ion flows across the cell membrane. Net ionic flux across the cell membrane determines the rate of change of V_m , with net inward flux causing cellular depolarization (increasing V_m) and net outward flux cellular repolarization (decreasing V_m). As a convention, all inward plasma membrane currents have negative polarity and ion transport across the cell membrane flows

into/out of S_1 . There are four plasma membrane conductances included in the UP model (Fig. 1), which are

1. inward Ca^{2+} current, I_{Ca} ,
2. plasma membrane Ca^{2+} -ATPase pump, I_{PM} ,
3. Ca^{2+} -inhibited nonselective cation conductance, I_{NSCC} ,
4. Outward Na^+ pump, I_{Na} .

For V_m near RMP, I_{Ca} provides the initial Ca^{2+} flux required for calcium-induced-calcium-release (CICR) initiation. Note that other conductances have also been identified in ICCs as being candidates for providing this initial Ca^{2+} influx; specifically, the dihydropyridine-sensitive L-Type and the dihydropyridine-resistant T-Type Ca^{2+} currents (8). However, the magnitudes of these currents are negligible near RMP (8). I_{Ca} is modeled by a passive current flowing against the Ca^{2+} electrochemical gradient, and takes the form

$$I_{Ca} = g_{Ca}(V_m - E_{Ca}). \quad (2)$$

We set g_{Ca} to 0.01 pS to provide sufficient Ca^{2+} entry for CICR initiation and to maintain resting C_{S1} levels. The Ca^{2+} Nernst potential, E_{Ca} , is given by the equation

$$E_{Ca} = \frac{RT}{2F} \log_e \left(\frac{C_O}{C_{S1}} \right), \quad (3)$$

where C_O is extracellular $[\text{Ca}^{2+}]$ of 1800 μM , T is body temperature of 310.16 K, R is the universal gas constant, and F is Faraday's constant.

The nonselective cation conductance is postulated to be the main inward current responsible for UP depolarization (11,32). Our model of I_{NSCC} conducts both Ca^{2+} and Na^+ ions and is given by the equation

$$I_{NSCC(Z)} = g_{NSCC(Z)}(V_m - E_{NSCC}), \quad (4)$$

where Z is either Na or Ca. Current-voltage relations show that the I_{NSCC} reversal potential, E_{NSCC} , is ~ 0 mV (33). The conductance of both ions through I_{NSCC} is modeled by the C_{S1} -inhibited Hill equation:

$$g_{NSCC(Z)} = \hat{g}_{NSCC(Z)} \left(\frac{K_{NSCC}^{h_{NSCC}}}{K_{NSCC}^{h_{NSCC}} + C_{S1}^{h_{NSCC}}} \right). \quad (5)$$

Voltage-ramp experiments performed on this conductance yielded $K_{NSCC} = 0.375 \mu\text{M}$ and $h_{NSCC} = 3$ (33), however, we used the value $K_{NSCC} = 0.12 \mu\text{M}$ in the model. This lower value was necessary to ensure appreciable activation of I_{NSCC} for changes in $[\text{Ca}^{2+}]$ around resting C_{S1} levels of 0.12 μM . The fitted values of $\hat{g}_{NSCC(Ca)}$ and $\hat{g}_{NSCC(Na)}$ were 0.12 pS and 220 pS, respectively. This is consistent with the assumption that Na^+ is the main charge carrier responsible for membrane electrical activity.

The plasma membrane Ca^{2+} -ATPase pump extrudes Ca^{2+} across the cell membrane to regulate cytosolic $[\text{Ca}^{2+}]$. We model I_{PM} by a second-order Hill equation that is activated by C_{S1} :

$$I_{PM} = g_{PM} \left(\frac{K_{PM}^2}{K_{PM}^2 + C_{S1}^2} \right). \quad (6)$$

The value $K_{PM} = 1.00 \mu\text{M}$ was based on typical values from pancreatic acinar cells (34) and we fit $g_{PM} = 420 \text{ fA}$ to control resting C_{S1} levels.

The outward Na^+ pump accounts for the UP repolarization phase by extruding Na^+ ions from the pacemaker unit. Note that UP repolarization is not actually described by the Sanders hypothesis, therefore, I_{Na} is a phenomenological addition to the modeling framework. However, justification for a repolarizing conductance in the form of I_{Na} is provided in the Discussion. We model I_{Na} by the N_{S1} -activated Hill equation as

$$I_{Na} = g_{Na} \left(\frac{N_{S1}^{h_{Na}}}{K_{Na}^{h_{Na}} + N_{S1}^{h_{Na}}} \right). \quad (7)$$

We set $K_{Na} = 1.0 \times 10^4 \mu\text{M}$ and $h_{Na} = 4$ to ensure I_{Na} was sensitive to changes in physiological N_{S1} levels. The value of g_{Na} was fit to a value of $1.5 \times 10^4 \text{ fA}$ to ensure sufficient extrusion of Na^+ from the pacemaker unit so as to prevent transient increases in N_{S1} levels. The differential equation governing membrane potential is given by the equation

$$\frac{dV_m}{dt} = -\frac{1}{C_m} (I_{Ca} + I_{PM} + I_{NSCC(Ca)} + I_{NSCC(Na)} + I_{Na}), \quad (8)$$

where C_m is the cellular membrane capacitance and was set to a value of 20 pF, which is within the range of capacitances recorded from murine ICCs (32).

The IP_3 receptor model

The IP_3R releases Ca^{2+} from ER stores into the cytosol (or in terms of the UP model, S_2). The IP_3R representation used with the modeling framework is based on a simplified two-state model (35). In modeling J_{IPR} , we use the same assumptions given in Sneyd et al. (35), namely, that the IP_3R is comprised of four identical and independent subunits and can only conduct Ca^{2+} if all the subunits are in an open state. In addition, it is also assumed that the IP_3R Ca^{2+} flux is proportional to the C_{ER}/C_{S2} gradient, giving the equation for J_{IPR} :

$$J_{IPR} = k_{IPR} \left(\frac{P\phi_1 H}{P\phi_1 + \phi_{-1}} \right)^4 (C_{ER} - C_{S2}), \quad (9)$$

where P is the $[\text{IP}_3]$, ϕ represents the IP_3R -state transitional rates, and H is a SV that dynamically controls the receptor's open state. All the equations governing H and ϕ are identical to the formulations given in Sneyd et al. (35) except for ϕ_3 , which we have modeled as a time-dependent “slow” SV. The purpose of this modification was to lengthen the quiescent phase of the oscillatory cycle by reducing the rate at which the IP_3R returns to a susceptible state. We deemed this modification necessary to reconcile the fast UP repolarization rates with the relatively slow pacing frequency ($1/Y_F \approx$

$100Y_{HW}$). We implemented a Hodgkin-Huxley type formulation to model ϕ_3 (36), which is given by the differential equation

$$\frac{d\phi_3}{dt} = \alpha_{\phi_3} - \beta_{\phi_3} \phi_3. \quad (10)$$

To obtain an initial feasible IP_3R parameter set that produced an oscillatory response, we used the parameter values quoted from Sneyd et al. (35). To increase the operable $[\text{Ca}^{2+}]$ range of the IP_3R , the parameters R_1 and R_3 were scaled by a factor of 6 to values of $36 \mu\text{M}$ and $300 \mu\text{M}$, respectively. All of the remaining IP_3R rate parameters were determined via the fitting procedure (see Parameter estimation) and are listed in Table 2. Note that we do not dynamically track P , because oscillations of this metabolite have not been demonstrated in ICCs (11). Therefore, we have assumed a constant value of $P = 1 \mu\text{M}$.

Other intracellular Ca^{2+} fluxes

The intracellular Ca^{2+} fluxes represent Ca^{2+} flows through the cytosolic volumes and the intracellular organelles. These fluxes are responsible for the intracellular Ca^{2+} oscillations and the localized Ca^{2+} depletion beneath the cell membrane required to activate I_{NSCC} . Our model includes five intracellular Ca^{2+} fluxes (Fig. 1):

1. Sarcoendoplasmic reticulum Ca^{2+} -ATPase, J_{SERCA} .
2. Mitochondrial Ca^{2+} uniporter, J_{MCU} .
3. Mitochondrial $\text{Na}^+/\text{Ca}^{2+}$ exchanger, J_{NCX} .
4. Intercytosolic subspace Ca^{2+} flux, J_{SIS2} .

The SERCA pump proteins transport Ca^{2+} ions from the cytosol into the ER, and are activated by increased cytosolic $[\text{Ca}^{2+}]$ (37,38) and/or by decreased luminal ER $[\text{Ca}^{2+}]$ (39). Our model of J_{SERCA} is based on the nonbuffering four-state model of Higgins et al. (40), which binds one Ca^{2+} ion and has the form

$$J_{SERCA} = \frac{V_{SERCA}(C_{S1} - A_2 C_{ER})}{1 + A_4 C_{S1} + A_5 C_{ER} + A_6 C_{S1} C_{ER}}. \quad (11)$$

The reason for this choice of SERCA pump model is because its formulation includes both cytosolic and luminal $[\text{Ca}^{2+}]$ dependencies. This is opposed to the majority of SERCA pump models, which only have cytosolic $[\text{Ca}^{2+}]$ dependencies (18,41) or have very weak luminal $[\text{Ca}^{2+}]$ dependencies (42). We used the parameter values $A_4 = 3.57 \mu\text{M}^{-1}$, $A_5 = 2.70 \times 10^{-5} \mu\text{M}^{-1}$, and $A_6 = 2.31 \times 10^{-5} \mu\text{M}^{-2}$ quoted from pancreatic acinar cells (40), and V_{SERCA} was fit to a value of $1 \times 10^5 \text{ s}^{-1}$. In determining the value of A_2 , we used the assumption that C_{ER} and C_{S1} reached near-steady-state levels and that the IP_3R has negligible Ca^{2+} flux during the quiescent phase of the UP cycle. From this assumption, we can approximate A_2 using the equation

$$A_2 = C_{S1(SS)}/C_{ER(SS)}, \quad (12)$$

where $C_{S1(SS)}$ is steady-state C_{S1} of $0.12 \mu\text{M}$ and $C_{ER(SS)}$ is steady-state C_{ER} , which we set to $\sim 200 \mu\text{M}$. Substituting these values into Eq. 12, we therefore obtained a value of $A_2 = 0.12/200 = 6.0 \times 10^{-4}$.

The MCU is a high-capacity, low- Ca^{2+} -affinity conductance that rapidly accumulates Ca^{2+} from the cytosol (or, in terms of the UP model, S_2). The J_{MCU} model is based on the formulation given in Gunter and Pfeiffer (30), which models the MCU by a second-order Hill equation as follows:

$$J_{MCU} = V_{MCU} \left(\frac{C_{S2}^2}{K_{MCU}^2 + C_{S2}^2} \right) \varepsilon_{INH}. \quad (13)$$

We set K_{MCU} at $10 \mu\text{M}$, which is within the range of values quoted from the literature for this conductance (30), and the value of V_{MCU} was set to $800 \mu\text{M s}^{-1}$ to ensure sufficient Ca^{2+} removal from S_2 . The MCU inhibition term, ε_{INH} is a phenomenological modification to prevent excessive build-up in mitochondrial Ca^{2+} levels over multiple cycles, and is given by the equation

$$\varepsilon_{INH} = \frac{K_{INH}^{h_{INH}}}{K_{INH}^{h_{INH}} + C_{MT}^{h_{INH}}}. \quad (14)$$

We used values of $K_{INH} = 10 \mu\text{M}$ and $h_{INH} = 4$ such that MCU activity was unaffected for normal (physiological) C_{MT} oscillations, but was significantly reduced for high (pathophysiological) C_{MT} levels.

The mitochondrial NCX is the primary mechanism of Ca^{2+} release from the mitochondrial matrix (30,43). We model this conductance using the formulation given in Colegrove et al. (44), which has the functional form

$$J_{NCX} = V_{NCX} \left(\frac{C_{MT}}{K_{NCX} + C_{MT}} \right). \quad (15)$$

We used a value of $K_{NCX} = 0.3 \mu\text{M}$, which is in agreement with estimates obtained from sympathetic neurons (44). A value of $V_{NCX} = 0.5 \mu\text{M s}^{-1}$ was used to prevent a transient build-up in mitochondrial Ca^{2+} over subsequent UP cycles. Note that the model of J_{NCX} only conducts Ca^{2+} because mitochondrial $[\text{Na}^+]$ changes are typically small and, as such, have a negligible effect on conductance activity (45).

The intercytoplasmic subspace Ca^{2+} flux provides a mechanism for Ca^{2+} diffusion between the cytoplasmic volumes. We use a simplification of the general diffusion equation (46) to model J_{S1S2} as follows:

$$J_{S1S2} = \mu_{S1S2} (C_{S2} - C_{S1}). \quad (16)$$

We used a value of $\mu_{S1S2} = 0.04 \text{ s}^{-1}$ to 1), allow sufficient Ca^{2+} entry from S_1 to initiate CICR, and 2), prevent Ca^{2+} diffusing back into the bulk cytosol before the MCU could properly activate.

Mitochondrial Ca^{2+} buffering

The mitochondrial Ca^{2+} buffer is a phenomenological addition to the model that is used to prevent depletion of C_{MT}

stores. The Ca^{2+} buffer model we used is based on the fast-approximation Ca^{2+} buffer model (47), and is given by the equation

$$f_m = \frac{1}{1 + \frac{K_m B_m}{(K_m + C_{MT})^2}}. \quad (17)$$

Note that Eq. 17 differs from the original formulation in that f_m appears on the righthand, as opposed to the lefthand, side of the differential equation. The buffering parameters, K_m and B_m , were set to values of $0.01 \mu\text{M}$ and $100 \mu\text{M}$, respectively such that prolonged net Ca^{2+} efflux from the mitochondria did not fully deplete Ca^{2+} stores.

Pacemaker unit model

Now that all of the constituent equations have been defined, the differential equations governing the intracellular ionic species can be stated. With reference to Fig. 1, the ionic species conservation ordinary differential equations are

$$\frac{dC_{S1}}{dt} = (J_{S1S2} + \lambda_{MT/S1} J_{NCX}) - ((\delta_S/Z_{Ca}) I_{Ca} + \lambda_{ER/S1} J_{SERCA}); \quad (18)$$

$$\frac{dC_{S2}}{dt} = \lambda_{ER/S2} J_{IPR} - (\lambda_{S1/S2} J_{S1S2} + \lambda_{MT/S2} J_{MCU}); \quad (19)$$

$$\frac{dC_{ER}}{dt} = J_{SERCA} - J_{IPR}; \quad (20)$$

$$\frac{dC_{MT}}{dt} = f_m (J_{MT} - J_{NCX}); \quad (21)$$

$$\frac{dN_{S1}}{dt} = -(\delta_S/Z_{Na}) I_{Na}, \quad (22)$$

where δ_S is a scale factor that converts a plasma membrane current, in pA, to an intracellular flux, in $\mu\text{M s}^{-1}$, and where Z_X is the valency of ion X . Note that for convenience, the plasma membrane conductances have been divided into separate Na^+ and Ca^{2+} aggregate currents:

$$I_{Ca} = I_{Ca} + I_{NSCC(Ca)} + I_{PM}; \quad (23)$$

$$I_{Na} = I_{NSCC(Na)} + I_{Na}. \quad (24)$$

The intercompartmental volume ratio, $\lambda_{X/Y}$, which scales the ionic fluxes between compartmental volumes, hence conserving ionic concentrations, is calculated as

$$\lambda_{X/Y} = \gamma_X / \gamma_Y, \quad (25)$$

where X and Y are the compartmental volumes.

A complete list of all the UP model equations is given in Appendix A, and all model parameters and SV initial conditions are given in Tables 2 and 3. Numerical simulations were carried out on a 3.0 GHz Intel Pentium 4 desktop computer using MATLAB (Version 7.1.0.246 Release 14; The MathWorks, Natick, MA). The model differential equations were integrated using the stiff equation solver, ode15s. The nondefault ode15s integration options were AbsTol = 1.0E-4 and RelTol = 1.0E-8.

TABLE 2 UP model parameters

	Parameter	Parameter value	Equations	Reference
Plasma membrane current parameters				
	g_{Ca}	0.01 pS	A12	F
	E_{NSCC}	0 mV	A14	(33)
	$\hat{g}_{NSCC(Ca)}$	0.12 pS	A15	F
	$\hat{g}_{NSCC(Na)}$	220 pS	A15	F
	K_{NSCC}	0.12 μ M	A15	F
	h_{NSCC}	3	A15	(33)
	g_{PM}	420 fA	A16	F
	K_{PM}	1 μ M	A16	(34)
	g_{Na}	1.5×10^4 fA	A17	F
	K_{Na}	1×10^4 μ M	A17	F
	h_{Na}	4	A17	F
Intracellular Ca^{2+} flux parameters				
	V_{SERCA}	1×10^5 s ⁻¹	A18	F
	A_2	6×10^{-4}	A18	F
	A_4	3.57μ M ⁻¹	A18	(40)
	A_5	$2.7 \times 10^{-5} \mu$ M ⁻¹	A18	(40)
	A_6	$2.31 \times 10^{-5} \mu$ M ⁻²	A19	(40)
	V_{MCU}	800μ M s ⁻¹	A19	F
	K_{MCU}	10 μ M	A19	(30)
	K_{INH}	10 μ M	A20	F
	h_{INH}	4	A20	F
	V_{NCX}	0.5μ M s ⁻¹	A21	F
	K_{NCX}	0.3 μ M	A21	(44)
	μ_{SIS2}	0.04 s ⁻¹	A22	F
IP ₃ R parameters				
	k_{IPR}	2000 s ⁻¹	A23	F
	k_1	0 s ⁻¹	A24	F
	k_{-1}	6.4μ M s ⁻¹	A25	F
	k_2	4 s ⁻¹	A26	F
	r_2	200 s ⁻¹	A24	F
	r_{-2}	0μ M s ⁻¹	A25	F
	r_4	750 s ⁻¹	A26	F
	R_1	36 μ M	A24	(35)
	R_3	300 μ M	A25,26	(35)
	g_{α}	0.02 s ⁻²	A28	F
	g_{β}	300 s ⁻¹	A29	F
	K_{β}	2 μ M	A29	F
	h_{β}	2	A29	F
Mitochondrial Ca^{2+} buffer parameters				
	K_m	0.01 μ M	A10	F
	B_m	100 μ M	A10	F
Intracellular subspace relative volumes				
	γ_{S1}	100	A11	(12,11)
	γ_{S2}	1	A11	(12,11)
	γ_{ER}	20	A11	(12,11)
	γ_{MT}	200	A11	(12,11)
Other constants				
	δ_S	26 μ M C ⁻¹	A1,6	F
	C_m	20 pF	A1	(32)
Cell constants				
	Z_{Ca}	2	A2	
	Z_{Na}	1	A6	
	R	8.31×10^{-3} aJ zmol ⁻¹ K ⁻¹	A13	
	T	310.16 K	A13	
	F	0.09649 fC zmol ⁻¹	A13	
	C_O	1.8×10^3 μ M	A13	
	P	1 μ M	A23	

F, parameter value was fitted.

TABLE 3 State variable initial conditions

State variable	Initial value
V_m	-70.1 mV
C_{S1}	0.120 μ M
C_{S2}	0.023 μ M
C_{ER}	203 μ M
C_{MT}	0.220 μ M
N_{S1}	1.01×10^4 μ M
ϕ_3	0.306 s ⁻¹

RESULTS

Unitary potential depolarizations

Integration of the model equations produces UPs with amplitudes of 3.00 mV (from a minimum diastolic potential of -70.2 mV), frequency of 0.05 Hz, and half-width of 0.21 s (Fig. 3 A). These UPs are spontaneously rhythmic (Fig. 3 B), and long-term simulations of more than 1000 oscillations show that the depolarizations are stable (data not shown). Table 4 summarizes all the model UP characteristic values and compares them to their experimentally recorded analogous values (9).

Intracellular ion cycling

Underlying the UP depolarizations are oscillations in the intracellular ionic species (Fig. 4). The model oscillations begin with Ca^{2+} entry across the cell membrane, which passively diffuses from S_1 to S_2 , transiently increasing both C_{S2} and the IP₃R open probability (Fig. 5 A). Once the IP₃R open probability reaches CICR threshold, ER Ca^{2+} is released into the microdomain, causing a significant increase in C_{S2} (Fig. 5, B and C). This has the effect of simultaneously inactivating the IP₃R and activating the MCU, causing rapid

mitochondrial Ca^{2+} accumulation (Fig. 5 D). Note that at this point the bulk of the Ca^{2+} released from the ER is now located in the mitochondria. Therefore, when the SERCA pumps sequester Ca^{2+} from S_1 to refill ER stores, C_{S1} falls below its original levels, causing I_{NSCC} activation (Fig. 5 E). The resultant net ionic influx depolarizes the cell membrane, with Ca^{2+} entry restoring C_{S1} levels, which inactivates I_{NSCC} , and Na^+ entry elevating N_{S1} , activating I_{Na} (Fig. 5 F). Therefore, the simultaneous inactivation of I_{NSCC} and activation of I_{Na} shifts net transmembrane ion flow from inward to outward, hence repolarizing the cell. Restoration of the intracellular ionic species returns the pacemaker unit to a susceptible state, allowing this cycle to continue ad infinitum. The process described above is summarized by the schematic diagram given in Fig. 6.

Cytosolic Ca^{2+} oscillations

Cytosolic $[Ca^{2+}]$, calculated as the weighted average of C_{S1} and C_{S2} , produces a spike over the pacemaker cycle (Fig. 7). The increase in cytosolic $[Ca^{2+}]$ is attributed to the large C_{S2} transient, which offsets the decrease in C_{S1} . This is despite the fact that C_{S1} makes up the bulk of cytosolic $[Ca^{2+}]$. This result is qualitatively similar to spatially averaged ICC intracellular Ca^{2+} image recordings (48–52). An interesting point to note about the simulated cytosolic $[Ca^{2+}]$ traces is that despite the reduction in subplasma membrane, required to activate pacemaker currents, there is still a general increase in $[Ca^{2+}]$ over the pacemaker cycle.

Voltage-clamp simulations

Simulations of voltage-clamp experiments were performed to determine the effect that membrane potential has on pace-

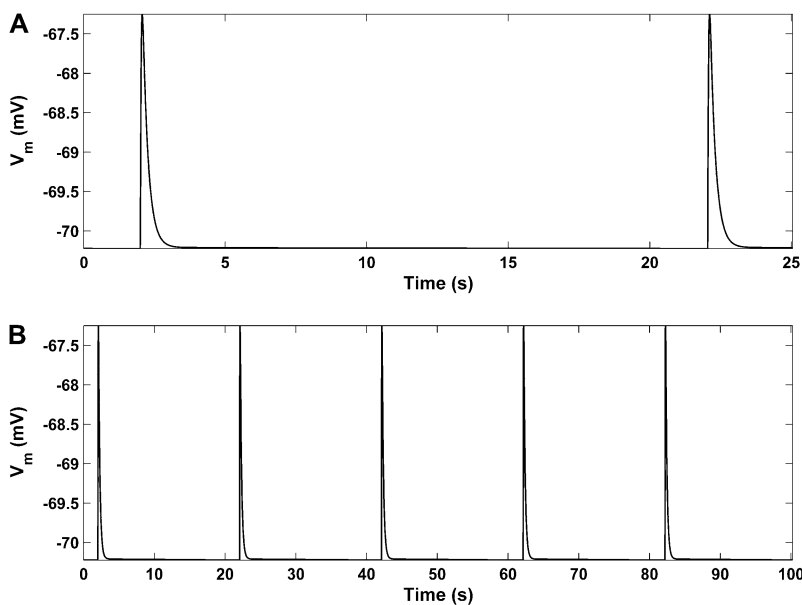


FIGURE 3 UP membrane depolarizations produced from model simulations. (A) A single UP oscillation. (B) UPs over multiple pacemaker cycles, illustrating the spontaneous rhythmicity of model pacemaker activity.

TABLE 4 Comparison of simulated and experimentally recorded UP characteristic values

Characteristic value	UP model	Experiment recordings
Amplitude (mV)	3.00	3.00*
Frequency (Hz)	0.05	0.05
Half-width (s)	0.21	0.21 [†]

*Modal amplitude.

[†]Values estimated from Fig. 2 *B* of Kito et al. (9).

maker activity. To replicate voltage-clamp experimental conditions, simulations were run with Eq. 8 set to zero and V_m set to the clamp potential. All other model state variables were calculated normally. The clamp potentials used for these simulations were 1) -70 mV to simulate control activity; 2), -60 mV to simulate depolarized conditions; and 3), -80 mV to simulate hyperpolarized conditions.

Results of these simulations qualitatively demonstrate two experimentally observed phenomena. First, model pacemaker activity persists under voltage clamp, as signified by the continuation of intracellular ion oscillations (Fig. 8). This result is comparable to experimental studies wherein voltage-clamped ICCs continued to generate pacemaker currents (25,32,33,53). Second, the reduction in cytosolic $[Ca^{2+}]$ caused by membrane depolarization (Fig. 8 *E*) reduces pacemaker frequency. This result is consistent with experimental observations, made by multiple investigators, that reduction in ICC intracellular $[Ca^{2+}]$ induced by the Ca^{2+} chelator, BAPTA-AM, caused a reduction in UP discharge rate (7,9,10).

Effects of blocking mitochondrial Ca^{2+} uptake

To determine the mitochondria's influence on pacemaker activity, numerical experiments were performed in which mitochondrial Ca^{2+} uptake was blocked. To replicate the experimental conditions of Beckett et al. (10), MCU activity was negated by collapsing the mitochondrial Ca^{2+} electrochemical gradient (i.e., setting Eq. 11 to zero). No alterations were made to any of the other model parameters.

The results of this numerical investigation show that inhibition of the MCU abolishes spontaneous membrane depolarizations (Fig. 9 *A*). The cessation of pacemaker activity appears to be caused by a build-up in C_{S2} that prevents the IP_3R from returning to a fully activated state (Fig. 9, *B* and *C*). Prolonged simulation of the model, after the cessation of pacemaker activity, shows that there is significant membrane hyperpolarization caused by a transient increase in C_{S1} (Fig. 9, *D* and *E*). This increase in C_{S1} is due to the emptying of mitochondrial Ca^{2+} stores over time (Fig. 9 *F*). Both of these results are similar to observations from experiments on murine gastric fundus, wherein mitochondrial Ca^{2+} uptake had been inhibited (10).

DISCUSSION

In this study, we introduce the first, to our knowledge, biophysically based computational modeling framework that simulates ICC pacemaker activity on the pacemaker unit spatial scale. This modeling framework is also the first to take into account a variety of biophysical parameters that have

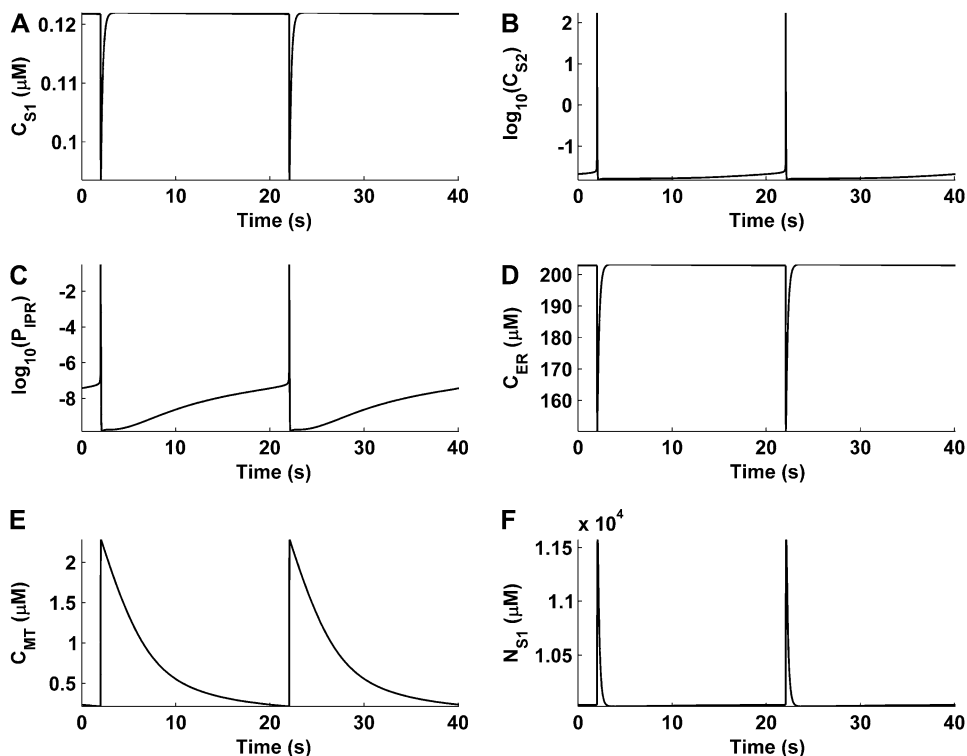


FIGURE 4 Intracellular ionic concentrations and IP_3R open probability tracked over multiple pacemaker cycles. The plots displayed are (A) S_1 $[Ca^{2+}]$, C_{S1} , (B) S_2 $[Ca^{2+}]$, C_{S2} , (C) IP_3R open probability, P_{IPR} , (D) ER $[Ca^{2+}]$, C_{ER} , (E) mitochondrial $[Ca^{2+}]$, C_{MT} , and (F) S_1 $[Na^+]$, N_{S1} . Note that logarithmic values are plotted for the C_{S2} (B) and P_{IPR} (C) traces.

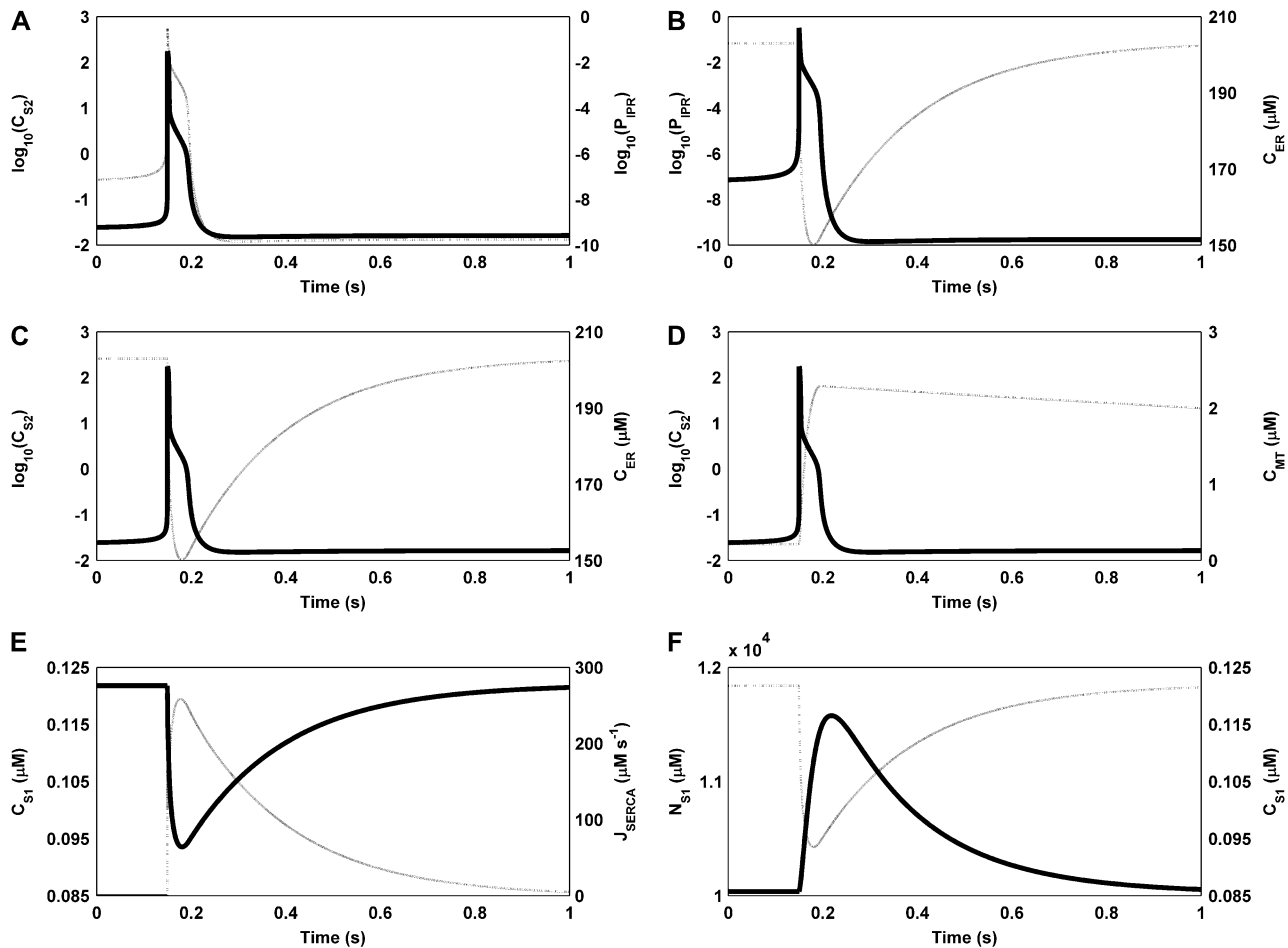


FIGURE 5 Plots showing the temporal relationship between the model SVs during the ER Ca^{2+} release phase of the UP oscillatory cycle. (A) The simultaneous increase in C_{S2} and P_{IPR} leads to (B) initiation of ER Ca^{2+} causing a sudden reduction in C_{ER} , which results in (C) a significant increase in C_{S2} . (D) The significant increase in C_{S2} activates the MCU, causing mitochondrial Ca^{2+} accumulation. (E) Simultaneously, the SERCA pumps activate to replenish ER Ca^{2+} stores, hence causing a reduction in C_{S1} . The depletion in C_{S1} activates I_{NSCC} , resulting in (F) a significant influx of Na^+ into the pacemaker unit. Note that the SV on the left vertical axis is represented by the solid curve, and the SV on the right vertical axis by the dashed curve. Also note that logarithmic values are plotted for the C_{S2} (B) and P_{IPR} (C) traces.

been measured experimentally: 1), the involvement of mitochondria as a basic mechanism in generating spontaneous inward currents; 2), provisions for a transient reduction in subplasma membrane $[\text{Ca}^{2+}]$, which regulates the open probability of a Ca^{2+} -inhibited nonselective cation conductance expressed by pacemaker ICCs; 3), modeling of Ca^{2+} dynamics within the constraints of an anatomical ICC ‘‘pacemaker unit’’. Previous models, though capable of generating rhythmical ICC slow-wave-like depolarizations, have not included any of these elements. Furthermore, all of the established models, except that of Edwards and Hirst (12), assume that slow waves (or pacemaker potentials), and not UP depolarizations, are the basic ICC pacemaker event.

Our modeling framework was based on the Sanders pacemaker hypothesis (11,15), which proposes a complex intracellular Ca^{2+} handling mechanism to activate pacemaker currents. In constructing our model, a minimalist approach was used wherein components were included only if there was suitable justification. This method is juxtaposed against

the approach of altering established mathematical cellular models to fit a desired framework. The risk of altering established models is that there could be intrinsic mechanisms that are neither necessary nor justifiable in the new framework. Consequently, by constructing the model from an elementary basis that is consistent with the biophysical mechanisms of the Sanders hypothesis, we have produced a model that is, we believe, the simplest and most accurate biophysical representation required to produce ICC subcellular pacemaker activity.

Here, we discuss the model’s performance, validation simulation results, a comparison of our model to the original hypothesis, the results of the model parameter sensitivity analysis, and the model’s limitations.

Physiological simulations

Results of the physiological numerical simulations show that the important aspects of pacemaker activity are depicted well by our model. Perhaps the most essential feature is the

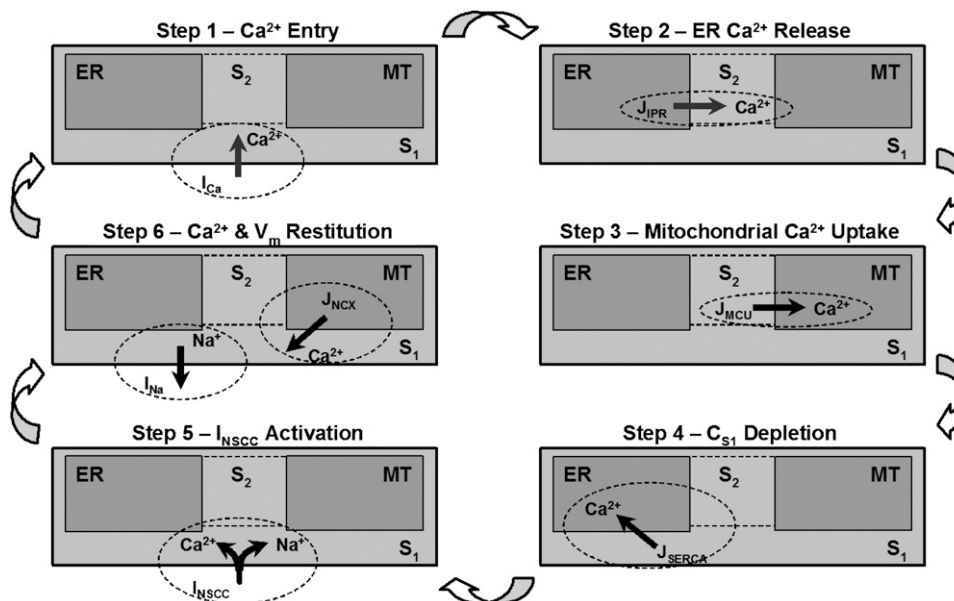


FIGURE 6 Schematic diagram of the pacemaking mechanism cycle as illustrated by the model response. (Step 1) Initial Ca^{2+} entry from the V_m -dependent inward Ca^{2+} current, I_{Ca} , which diffuses from S_1 to S_2 . (Step 2) Sufficient Ca^{2+} entry into S_2 raises the IP_3R open probability to threshold and ER Ca^{2+} release occurs. (Step 3) The subsequent increase in C_{S_2} activates the MCU, causing rapid mitochondrial Ca^{2+} accumulation. (Step 4) ER Ca^{2+} sequestration via the SERCA pumps to replenish ER Ca^{2+} stores. (Step 5) Activation of I_{NSCC} , caused by C_{S_1} depletion from SERCA activation. The resulting Ca^{2+} and Na^+ influx is responsible for the UP depolarization phase. (Step 6) Repolarization, caused by I_{Na} activation from increased N_{S_1} levels, and C_{MT} restitution, from J_{NCX} , resets the pacemaker unit allowing the cycle to begin again.

model's ability to produce spontaneously rhythmic membrane depolarizations (Fig. 3 B). This demonstrates that our model is capable of producing pacemaker activity, a property intrinsic to ICCs (11). Moreover, the UPs produced by our model have characteristics similar to those recorded experimentally (Table 4). Note that it is not actually possible to make a direct comparison between model simulations and experimentally recorded UP traces. This is because 1), experimentally recorded UPs are highly stochastic in nature, and 2), all experimentally recorded UPs are made from whole-cell or multiple-cell preparations, making it impossible to isolate the electrical behavior of a single pacemaker unit.

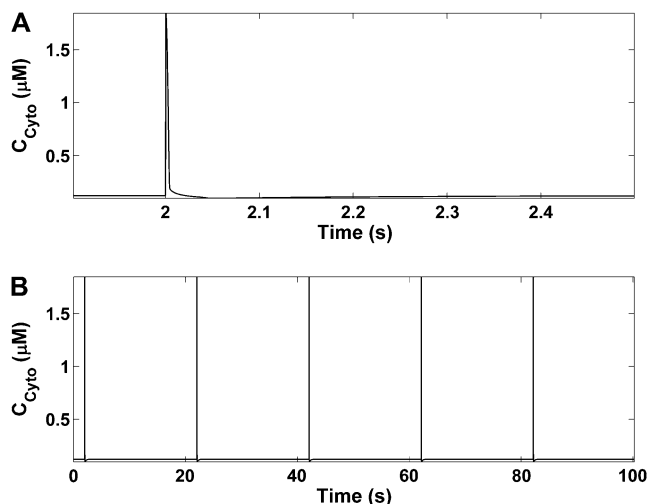


FIGURE 7 Pacemaker unit cytosolic $[\text{Ca}^{2+}]$ oscillations around the ER Ca^{2+} release phase of the pacemaker cycle (A) and over multiple pacemaker cycles (B). Note that there is an increase in cytosolic $[\text{Ca}^{2+}]$ over each cycle. This is despite the fact that depletion of C_{S_1} is required for the activation of pacemaker currents (I_{NSCC}) and C_{S_1} makes up the bulk of cytosolic $[\text{Ca}^{2+}]$.

Another important feature of our model is that the activation of pacemaker currents is caused by depletion of subplasma membrane $[\text{Ca}^{2+}]$. This result is important, because membrane depolarizations caused by the activation of a Ca^{2+} -inhibited conductance have not been demonstrated previously, that we know of, in ICC mathematical models. Instead, investigators have invariably relied upon a Ca^{2+} -activated conductance as a mechanism for generating transient inward currents and membrane depolarization, but there is no direct experimental evidence demonstrating the existence of a Ca^{2+} -activated inward conductance in pacemaker ICCs (11). A Ca^{2+} -inhibited nonselective cation conductance has been identified and characterized by ICC experiments (33).

Analysis of the pacemaker unit cytosolic $[\text{Ca}^{2+}]$ shows that there is an increase in $[\text{Ca}^{2+}]$ over the oscillatory cycle (Fig. 7). This result is comparable to spatially averaged images of intracellular $[\text{Ca}^{2+}]$ in ICC (48–52). Note that a direct comparison cannot be made between the experimentally recorded and numerically simulated Ca^{2+} traces. The reasons for this are that 1), the experimental recordings are a relative, not absolute, measure of intracellular $[\text{Ca}^{2+}]$, and 2), our mathematical model is only representative of pacemaker activity from a single pacemaker unit, whereas the experimental recordings represent whole-cell activity, where voltage-dependent Ca^{2+} currents may amplify cytosolic transients. Nonetheless, this result overcomes a seemingly paradoxical situation: the depletion of subplasma membrane Ca^{2+} required for I_{NSCC} activation while still producing an increase in cytosolic $[\text{Ca}^{2+}]$ over the oscillatory cycle, which is qualitatively similar to experimental observations.

Finally, our model demonstrates that ICC pacemaker activity is possible in the absence of IP_3 oscillations. This result is important, because, despite the consensus that release of

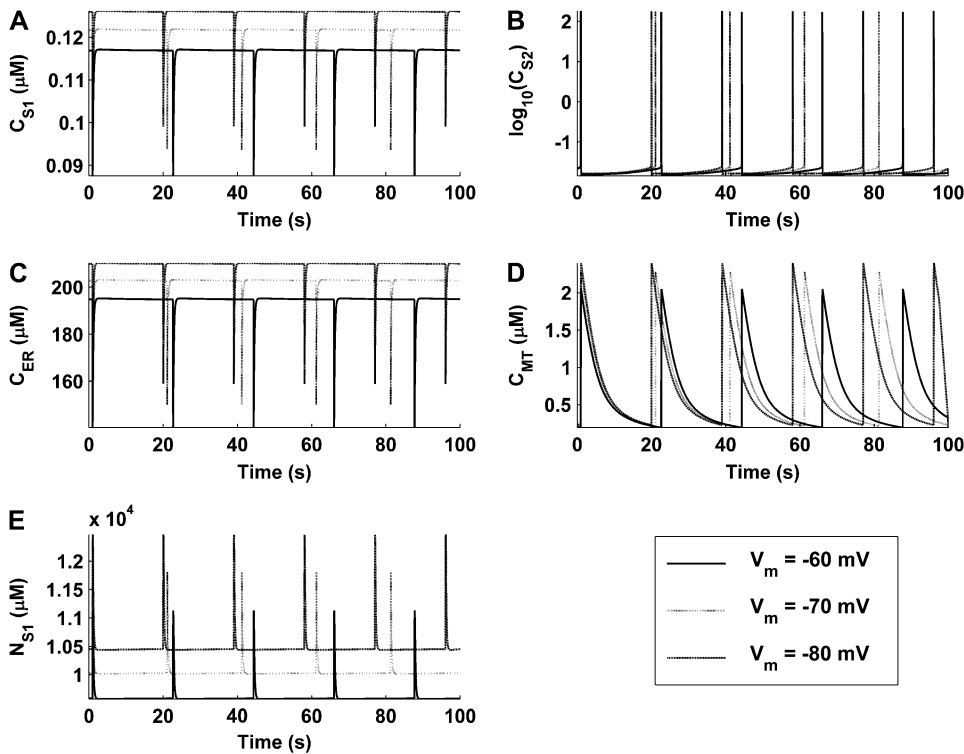


FIGURE 8 Intracellular ionic concentrations produced under simulated voltage-clamp conditions for (A) C_{S1} , (B) C_{S2} , (C) C_{ER} , (D) C_{MT} , and (E) N_{S1} . Simulations were performed for $V_m = -80$, -70 , and -60 mV. The simulations show that depolarization causes a decrease in pacemaker frequency and a decrease in the magnitude of the intracellular ionic concentrations (the effect is opposite for hyperpolarization). Note that logarithmic values are plotted for the C_{S2} trace (B).

Ca^{2+} via IP_3 -operated ER stores is essential for ICC pacemaker activity (11,14,25), there remains controversy over the exact mechanism by which IP_3 R-operated Ca^{2+} release is initiated. Several investigators have suggested that V_m -dependent IP_3 synthesis is responsible for regulating ER Ca^{2+} release (25), and indeed such a mechanism has been utilized in other ICC mathematical models (13,20). In our modeling framework, Ca^{2+} entry initiates ER Ca^{2+} release. The problem with the V_m -dependent IP_3 synthesis mechanism is that $[IP_3]$ oscillations phase-locked to pacemaker activity have never been demonstrated experimentally in ICCs (11). This is coupled with the fact that the predominant IP_3 R isoform in ICCs is the Type-1 isoform (54), which is characterized by a high- Ca^{2+} -gain, low- IP_3 -gain response (55). As a consequence, the ICC IP_3 R is more sensitive to changes in $[Ca^{2+}]$ in comparison to changes in $[IP_3]$. Conversely, pacemaker frequency and intracellular $[Ca^{2+}]$ oscillations in ICCs are clearly dependent upon Ca^{2+} entry mechanisms (48–52), giving credibility to the Ca^{2+} entry hypothesis. Moreover, pacemaker activity simulated by our model is further support for the involvement of a Ca^{2+} entry mechanism.

Validation simulations

Voltage-clamp experiments were performed on our model to investigate the dependency of simulated pacemaker activity on membrane potential. Results of these simulations predict that pacemaker activity persists under voltage clamp (Fig. 8),

a phenomenon that has been observed experimentally by multiple investigators (25,32,33,50). Simulations performed under increasingly depolarized clamp potentials show that there is a reduction in pacemaker frequency (Fig. 8). This is due to the overall reduction in the cytosolic Ca^{2+} load caused by a reduction in the magnitude of the inward Ca^{2+} currents, I_{Ca} and $I_{NSCC(Ca)}$ (Fig. 8 E). Indeed, pharmacological agents that reduce intracellular Ca^{2+} have been shown to decrease UP activity in ICCs (7,9,10). However, experimental observations have also shown that membrane depolarization increases UP activity (9,56). Therefore, as model pacemaker frequency is dependent on cytosolic Ca^{2+} load, which in turn is dependent on V_m -dependent Ca^{2+} entry, these results suggest the possibility that the modeling framework, in its present form, is “missing” an inward V_m -dependent Ca^{2+} conductance. This is not entirely unexpected because, as previously stated, the present modeling framework is only representative of pacemaker activity for V_m near RMP (see also Limitations, below).

Numerical investigations were also undertaken to determine the influence that the mitochondria has on pacemaker activity. To determine this influence, we removed the effect of MCU Ca^{2+} uptake by collapsing the mitochondrial Ca^{2+} electrochemical gradient. Results of these simulations show that pacemaker activity is abolished, with prolonged simulation predicting hyperpolarization of the cell membrane. These results are consistent with experimental observations in the guinea pig gastric fundus, where mitochondrial Ca^{2+} uptake was inhibited (10).

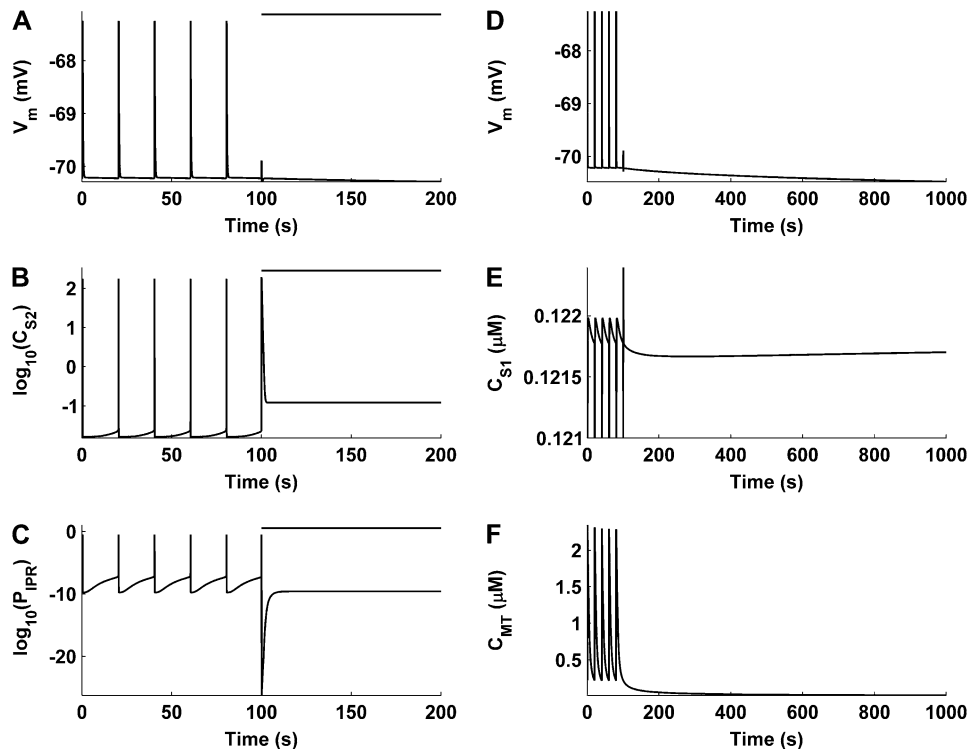


FIGURE 9 Numerical simulations showing the short-term (A–C) and long-term (D–F) effects of mitochondrial inhibition on pacemaker activity. The bar above the traces in A–C denotes MCU inhibition. UP depolarizations, which are present before MCU inhibition, are abolished after MCU block (A). This appears to be due to a build-up in C_{S2} (B) which prevents the IP_3R from returning to a susceptible state (C). Long-term simulation, after pacemaker activity cessation, shows membrane hyperpolarization (D) caused by a transient increase in C_{S1} (E). The increase in C_{S1} is due to the emptying of mitochondrial Ca^{2+} stores over time (F). Note that logarithmic values are plotted for the C_{S2} (B) and P_{IPR} (C) traces, and that the y axis on C_{S1} trace (E) is truncated at $0.121 \mu M$.

Hypothesis comparison

Our UP model has a biophysical basis that encapsulates the central mechanisms postulated by the Sanders pacemaker hypothesis. These mechanisms, as discussed above, include the V_m -dependent entry of Ca^{2+} to initiate CICR, the importance of ER Ca^{2+} cycling to pacemaker activity, the inclusion of mitochondrial Ca^{2+} dynamics, and the localized depletion of subplasma membrane $[Ca^{2+}]$ to initiate pacemaker currents. However, there are several points of difference between the UP model and the Sanders pacemaking hypothesis. These differences include the addition of a second cytosolic subspace, the mechanism by which subplasma membrane $[Ca^{2+}]$ is depleted, and the addition of a repolarizing current in the form of an outward Na^+ current. Each of these points of difference is discussed in turn below.

Addition of second cytosolic subspace

To account for the high Ca^{2+} activity in the region formed between the ER and mitochondria, a secondary cytosolic subspace was added to the modeling framework. Via this structure, the assumption of ion concentration homogeneity within each compartmental volume remains valid. As discussed above, this secondary cytosolic subspace is also consistent with the “hotspot theory” (29,30). This theory postulates that high $[Ca^{2+}]$ in the microdomain formed between the intracellular organelles is capable of activating the low- Ca^{2+} -affinity MCU, an event which is not possible under normal cytosolic $[Ca^{2+}]$ oscillations.

Mechanism for C_{S1} depletion

The original hypothesis suggests that mitochondrial Ca^{2+} uptake is responsible for removing Ca^{2+} from S_1 , whereas our model predicts that C_{S1} depletion occurs by SERCA pump activation. Although mitochondrial Ca^{2+} uptake is a requisite for pacemaker activity (10,25,57), within our modeling framework we were unable to define a consistent set of parameter values that produced C_{S1} depletion via the MCU alone. This is primarily due to the fact that the MCU draws its Ca^{2+} from S_2 and hence does not directly influence C_{S1} . The issues associated with this mechanism are further compounded by the fact that, 1), the rate of intercytosolic subspace Ca^{2+} diffusion is limited, and 2), S_2 is significantly smaller than S_1 and thus would have little influence on its activity. These issues cannot be circumvented because, as previously stated, S_2 Ca^{2+} accumulation is required for MCU activation.

If the MCU were to cause submembrane Ca^{2+} depletion, C_{S2} would have to be maintained at a lower level than C_{S1} for a prolonged period of time. This is so that Ca^{2+} could passively diffuse from S_1 back to S_2 . The difficulty with this proposed mechanism is that mitochondrial Ca^{2+} uptake is a negative feedback process whereby MCU Ca^{2+} uptake causes a reduction in C_{S2} , which in turn reduces the magnitude of the MCU Ca^{2+} flux (see Eq. 13). Furthermore, the MCU is a low- Ca^{2+} -affinity conductance and thus would have negligible magnitude at low $[Ca^{2+}]$. With all of these facts taken together, the MCU would not have the capacity to maintain the required low C_{S2} levels and hence could not possibly cause C_{S1} depletion.

In contrast, to consider the mechanism of C_{S1} depletion caused by the SERCA pump proteins we need to review the events directly after CICR. Ca^{2+} released into S_2 will activate the MCU, which rapidly accumulates Ca^{2+} in the mitochondrial matrix. Simultaneously, the ER will sequester Ca^{2+} from the cytoplasm (or, in terms of the model, S_1) to replenish Ca^{2+} stores. Note that at this point the bulk of the Ca^{2+} released from the ER now resides in the mitochondria. Therefore when Ca^{2+} is sequestered back into the ER there will be a net reduction in C_{S1} , hence activating the pacemaker currents. Numerical simulations performed on our model provide proof of concept to this mechanism (Fig. 4). Note that this new mechanism does not invalidate the original hypothesis, because mitochondrial Ca^{2+} uptake is still necessary for pacemaker activity, as exemplified by the MCU inhibition numerical simulations (Fig. 9).

I_{Na} as a repolarizing conductance

The argument for the inclusion of a repolarizing conductance, in the form of I_{Na} , is based on enforcing charge conservation within the model. Let I_{ion} denote the sum of all the plasma membrane conductances (i.e., I_{ion} equals the sum of Eqs. 23 and 24). Consider the closed-form equation used in Edwards and Hirst (12) to represent UP depolarizations,

$$V_m = H \left(e^{-t/A} - e^{-t/B} \right)^3, \quad (26)$$

where A and B are time constants, and H is a scale factor. By differentiating Eq. 26 with respect to time, substituting it into Eq. 8, and rearranging, we obtain a closed-form expression for the I_{ion} transient that is required to reproduce Eq. 26,

$$I_{ion} = -3HC_m \left(\frac{dU}{dt} U^2 \right), \quad (27)$$

where $U = e^{-t/A} - e^{-t/B}$, and

$$\frac{dU}{dt} = - \left(\frac{e^{-t/A}}{A} - \frac{e^{-t/B}}{B} \right).$$

Using typical UP time constants of $A = 0.434$ s and $B = 0.077$ s (12) and setting $3HC_m = 1$, it can be seen that there are two distinct phases of net transmembrane ionic flux, 1), net ion influx, and 2), net ion efflux (Fig. 10). Based on the Sanders hypothesis, we can assume that net ionic influx can be attributed to the activation of I_{NSCC} . The mechanism for net ionic efflux is unclear, but it is likely that the trigger is not associated with altered $[Ca^{2+}]$, because Ca^{2+} entry via I_{NSCC} would restore submembrane $[Ca^{2+}]$ to original levels. Therefore, the only other possible repolarization triggers are V_m , whereby depolarization could activate the outward current, or increased $[Na^+]$ levels, caused by Na^+ influx through I_{NSCC} .

Consider the case of V_m as the primary repolarization mechanism. The most likely candidate conductances for such a mechanism would be an outward V_m -dependent K^+ channel. Indeed, several different types of K^+ channels have been

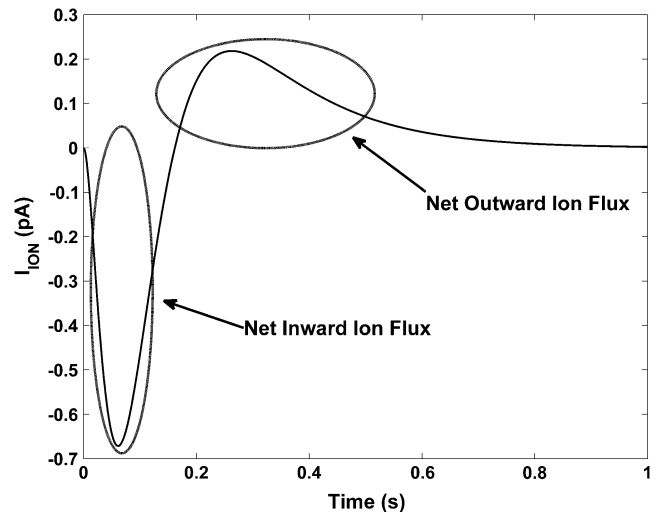


FIGURE 10 Trace of the aggregate plasma membrane current, I_{ion} , required to reproduce a UP depolarization using the UP representation of Edwards and Hirst (12) ($A = 0.434$ s, $B = 0.077$ s, $3HC_m = 1$). Note the two distinct phases of transmembrane ionic flux that comprise UP depolarizations: net inward ion flux and net outward ion flux.

identified in ICCs (32,58), and K^+ conductances are necessary for maintaining RMP in an excitable cell (59). It should be appreciated, however, that the amplitude of UP depolarizations are very small, typically in the range of 0.5–10 mV (7,9). This means that the voltage sensor of an outward V_m -dependent current would have to be extremely sensitive to these small depolarizations to produce the fast UP repolarization rates. However, the kinetics of the K^+ conductances identified from the aforementioned experimental studies show that these currents have negligible magnitude near RMP. Therefore, even though these conductances would eventually restore V_m depolarizations to original levels, the repolarization rate would not be fast enough to account for the UP half-width. As a result, K^+ conductances, although important for ICC pacemaker activity, are probably not directly involved in UP repolarization. Conversely, a repolarization mechanism dependent on elevated intracellular $[Na^+]$ is not constrained by the small amplitude of UP depolarizations. This is because Na^+ entry into the pacemaker unit is dependent only on the V_m -independent I_{NSCC} . Furthermore, an outward Na^+ conductance provides a pathway for Na^+ extrusion from the pacemaker unit. Taken altogether, we deemed a Na^+ -dependent outward conductance to be the most physiologically plausible option, and thus was the repolarization conductance implemented in the modeling framework.

We assume that the Na^+ -dependent outward conductance, denoted as I_{Na} , represents a Na^+ dependent ion exchanger such as a sodium/calcium or sodium/potassium exchanger. However as the exact identity of I_{Na} is unknown, we cannot be certain of the conductance's functionality or kinetics. Therefore, we conservatively model I_{Na} as a Na^+ -ATPase pump that has the same functionality and kinetics as I_{PM} (see

Eqs. 5 and 6), the only difference being that I_{Na} is dependent on intracellular $[Na^+]$ rather than $[Ca^{2+}]$.

Note that as I_{Na} is dependent on $[Na^+]$, the inclusion of S_1 $[Na^+]$, or N_{S1} , as a model SV, is justified. Note also that S_1 is the only compartmental volume where Na^+ dynamics are tracked by the model. This is because our primary concern is determining the effect of Na^+ on the plasma membrane, and all plasma membrane currents are regulated by S_1 ion concentrations.

Limitations

Although our model is capable of reproducing many important features of pacemaker activity, it must be appreciated that the model incorporates assumptions and features that are approximations of reality. Here, we identify and discuss the main limiting factors of the model.

The foremost limiting factor of ICC biophysically based modeling is the lack of experimental data describing dynamics such as ion channel kinetics, intracellular ion concentrations, and metabolic pathways. This limitation has also been noted by other biophysically based ICC modeling investigators (20). Due to the scarcity of suitable modeling data, a majority of our model parameter values were based on data from other cell types. For those parameters for which analogous values did not exist in the literature, we fit their values using optimization techniques to reproduce characteristic UP features. However, the biophysical structure of the model means that as experimental data becomes available we can readily adapt the model to incorporate this new information.

Another limitation of the model is that it is only representative of pacemaker activity for V_m near RMP. As a result, certain V_m -dependent conductances that are important for whole-cell pacemaker activity, but have negligible magnitude near RMP, have been omitted from the modeling framework. Such conductances include the dihydropyridine-resistant Ca^{2+} current, I_{VDDR} (8). Voltage-dependent Ca^{2+} entry via I_{VDDR} has been suggested to organize the discharge of unitary potentials from the pacemaker units (11,12), accounting for the upstroke phase of the PPs (16,17). Other important omissions are the K^+ channels, such as the Ba^{2+} -sensitive inwardly rectifying K^+ channels (32) and voltage-dependent $Kv1.1$ K^+ channels (58), which are suggested to account for the repolarization phase of the PPs (32). Therefore, numerical simulations that are performed in the voltage window of these omitted conductances ($\sim >60$ mV) do not take into account their activity, and hence these simulations are not an accurate representation of pacemaker activity.

SUMMARY

In this study, we introduce a mathematical model that is capable of reproducing ICC pacemaker activity on the sub-cellular spatial scale. Our goal for this project was to provide a biophysical basis for the complex mechanisms underlying GI pacemaker activity. This work is analogous to quantitative

electrophysiological studies that have been performed on other organs, such as the heart. Future work will allow us to exploit the multiscale approach of biophysically based mathematical modeling to entrain individual UPs to reproduce whole-cell slow-wave depolarizations. The ultimate goal is to embed an entire biophysically based ICC model into anatomically accurate GI whole-organ topologies to study electrical function under healthy and pathological conditions.

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APPENDIX A: MODEL EQUATION SUMMARY

The differential equations governing the model state variables are:

$$\frac{dV_m}{dt} = -\frac{1}{C_m}(I_{iCa} + I_{iNa}), \quad (A1)$$

$$\frac{dC_{S1}}{dt} = (J_{S1S2} + \lambda_{MT/S1}J_{NCX}) - ((\delta_S/Z_{Ca})I_{iCa} + \lambda_{ER/S1}J_{SERCA}), \quad (A2)$$

$$\frac{dC_{S2}}{dt} = \lambda_{ER/S2}J_{IPR} - (\lambda_{S1/S2}J_{S1S2} + \lambda_{MT/S2}J_{MCU}); \quad (A3)$$

$$\frac{dC_{ER}}{dt} = J_{SERCA} - J_{IPR}, \quad (A4)$$

$$\frac{dC_{MT}}{dt} = f_m(J_{MCU} - J_{NCX}), \quad (A5)$$

$$\frac{dN_{S1}}{dt} = -(\delta_S/Z_{Na})I_{iNa}, \quad (A6)$$

$$\frac{dH}{dt} = \phi_3(1 - H) - \frac{P\phi_1\phi_2}{P\phi_1 + \phi_{-1}}H, \quad (A7)$$

where

$$I_{iCa} = I_{Ca} + I_{NSCC(Ca)} + I_{PM}, \text{ aggregate } Ca^{2+} \text{ ionic current,} \quad (A8)$$

$$I_{iNa} = I_{NSCC(Na)} + I_{Na}, \text{ aggregate } Na^+ \text{ ionic current,} \quad (A9)$$

$$f_m = \frac{1}{1 + \frac{K_m B_m}{(K_m + C_{MT})^2}}, \text{ mitochondrial } Ca^{2+} \text{ buffering rate,} \quad (A10)$$

and

$$\lambda_{X/Y} = \frac{\gamma_X}{\gamma_Y}, \text{ compartmental volume ratio,} \quad (\text{A11})$$

where X and $Y = S_1, S_2, ER$, or MT .

The model plasma membrane currents are given as follows:

$$I_{Ca} = g_{Ca}(V_m - E_{Ca}), \text{ inward } Ca^{2+} \text{ leakage current;} \quad (\text{A12})$$

$$E_{Ca} = \frac{RT}{2F} \log_e \left(\frac{C_O}{C_{S1}} \right); \quad (\text{A13})$$

$$I_{NSCC(Z)} = g_{NSCC(Z)}(V_m - E_{NSCC}), \quad (\text{A14})$$

nonselective cation conductance;

and

$$g_{NSCC(Z)} = \hat{g}_{NSCC(Z)} \left(\frac{K_{NSCC}^{hNSCC}}{K_{NSCC}^{hNSCC} + C_{S1}^{hNSCC}} \right), \quad (\text{A15})$$

where $Z = Ca$ or Na and

$$I_{PM} = g_{PM} \left(\frac{C_{S1}^2}{K_{PM}^2 + C_{S1}^2} \right), \text{ plasma membrane } Ca^{2+}\text{-ATPase;} \quad (\text{A16})$$

and

$$I_{Na} = g_{Na} \left(\frac{N_{S1}^{hNa}}{K_{Na}^{hNa} + N_{S1}^{hNa}} \right), \text{ outward } Na^+ \text{ current.} \quad (\text{A17})$$

The model intracellular Ca^{2+} fluxes are

$$J_{SERCA} = \frac{V_{SERCA}(C_{S1} - A_2 C_{ER})}{1 + A_4 C_{S1} + A_5 C_{ER} + A_6 C_{S1} C_{ER}}, \text{ SERCA pump,} \quad (\text{A18})$$

$$J_{MCU} = V_{MCU} \left(\frac{C_{S2}^2}{K_{MCU}^2 + C_{S2}^2} \right) \varepsilon_{INH}, \quad (\text{A19})$$

mitochondrial Ca^{2+} uniporter,

$$\varepsilon_{INH} = \frac{K_{INH}^{hINH}}{K_{INH}^{hINH} + C_{MT}^{hINH}}, \quad (\text{A20})$$

$$J_{NCX} = V_{NCX} \left(\frac{C_{MT}}{K_{NCX} + C_{MT}} \right), \quad (\text{A21})$$

mitochondrial Na^+/Ca^{2+} exchanger,

$$J_{S1S2} = \mu_{S1S2}(C_{S2} - C_{S1}), \text{ intersubspace } Ca^{2+} \text{ flux,} \quad (\text{A22})$$

and

$$J_{IPR} = k_{IPR} \left(\frac{P\phi_1 H}{P\phi_1 + \phi_{-1}} \right)^4 (C_{ER} - C_{S2}), \text{ IP}_3\text{R } Ca^{2+} \text{ flux.} \quad (\text{A23})$$

The IP_3R rate constants are

$$\phi_1 = \frac{k_1 R_1 + r_2 C_{S2}}{R_1 + C_{S2}}, \quad (\text{A24})$$

$$\phi_{-1} = \frac{(k_{-1} + r_{-2}) R_3}{R_3 + C_{S2}}, \quad (\text{A25})$$

and

$$\phi_2 = \frac{k_2 R_3 + r_4 C_{S2}}{R_3 + C_{S2}}. \quad (\text{A26})$$

The IP_3R slow variable, ϕ_3 differential equation, and rate constants are

$$\frac{d\phi_3}{dt} = \alpha_{\phi_3} - \beta_{\phi_3} \phi_3; \quad (\text{A27})$$

$$\alpha_{\phi_3} = g_{\alpha}; \quad (\text{A28})$$

and

$$\beta_{\phi_3} = g_{\beta} \left(\frac{C_{S2}^{h\beta}}{K_{\beta}^{h\beta} + C_{S2}^{h\beta}} \right). \quad (\text{A29})$$

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